

# The effective Polyakov loop theory for finite temperature Yang-Mills theory and QCD

Georg Bergner  
ITP GU Frankfurt

O. Philipsen, J. Langelage, M. Neuman



Lattice: June 23, 2014

- 1 Effective Polyakov loop theory for QCD thermodynamics
- 2 Results at finite density from the effective theory
- 3 Low temperature heavy-dense regime
- 4 Numerical determination of effective couplings
- 5 Conclusions

## Effective Polyakov loop theory for QCD thermodynamics

$$\exp[-S_{\text{eff}}] \equiv \int [dU_k] \exp[-S_g(\beta)] \prod_{f=1}^{N_f} \det [Q^f(\kappa)]$$

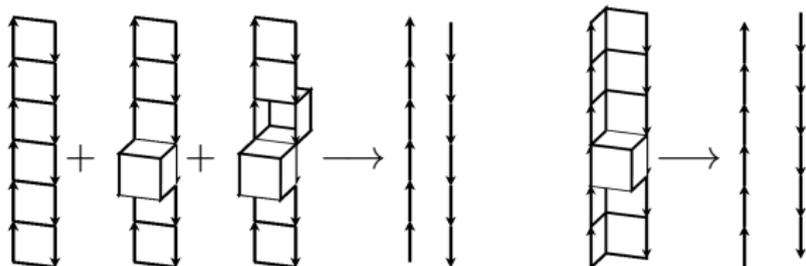
$$-S_g = \frac{\beta}{2N_c} \sum_p [\text{tr } U_p + \text{tr } U_p^\dagger]$$

- integration of spatial links  $U_k$ ; three dimensional theory  
 $U_\mu(x, t) \rightarrow U_0(x) \rightarrow$  Polyakov loops  $L(x)$

$$Z = \int [dL] e^{-S_{\text{eff}}[L]}$$

- $\Rightarrow$  simulations of effective theory allow to circumvent/reduce the sign problem at finite chemical potential

## The improved strong coupling approach



$$S_{\text{eff}} = \lambda_1 S_{\text{nearest neighbors}} + \lambda_2 S_{\text{next to nearest neighbors}} + \dots$$

- ordering in expansion parameter  $u = \frac{\beta}{18} + \dots < 1$
- simplest, most relevant contribution:

$$e^{-S_{\text{eff}}} \approx \prod_{\langle i,j \rangle \text{ nearest n.}} \left( 1 + 2\lambda_1 \Re(L_i L_j^\dagger) \right)$$

- $\lambda_1$ : resummed high orders in  $u$

## Hopping parameter expansion

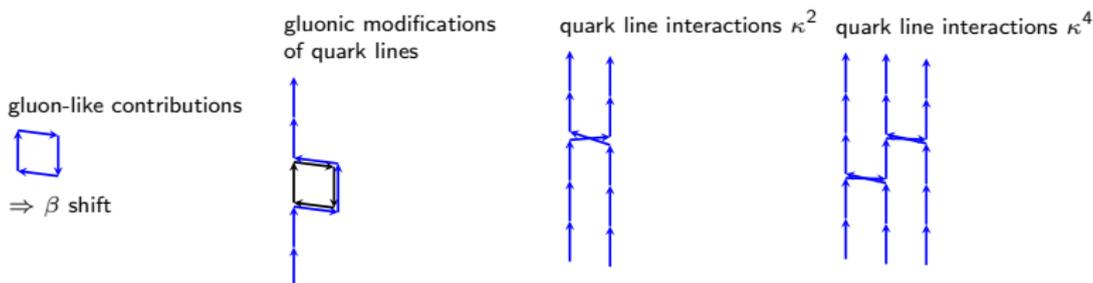
- Wilson-Dirac operator:  $Q = 1 - \kappa H[U]$
- $H = T + S$  with temporal  $T$  and spatial  $S$  hopping
- static contribution with resummation of windings

$$\begin{aligned} \det(1 - \kappa T) \\ = \prod_n (1 + cL_n + c^2L_n^\dagger + c^3)^2 (1 + \bar{c}L_n^\dagger + \bar{c}^2L_n + \bar{c}^3)^2 \end{aligned}$$

- chemical potential  $\mu$  in fugacity factor

$$c = (2\kappa e^{a\mu})^{N_\tau} = \exp\left(\frac{\mu - m_{stat}}{T}\right); \quad \bar{c} = (2\kappa e^{-a\mu})^{N_\tau}$$

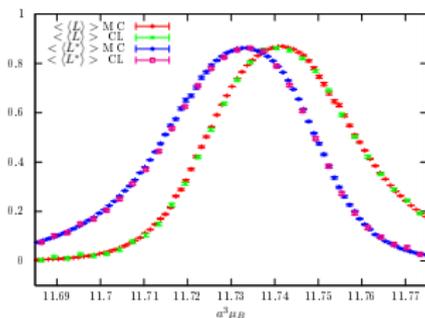
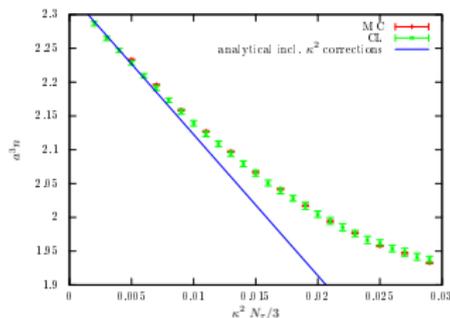
# Hopping parameter expansion: kinetic part



$$\det Q = \det(1 - \kappa T) \det(1 - (1 - \kappa T)^{-1} S)$$

- kinetic term  $\det(1 - (1 - \kappa T)^{-1} S)$ : expansion in  $\kappa$
  - gluonic modifications by plaquette contributions ( $u$ )
- $\Rightarrow$  interactions up to  $\kappa^n + u^m$ ,  $m + n = 4$  included
- $\Rightarrow \kappa^n$  contributions can be automatized for higher  $n$

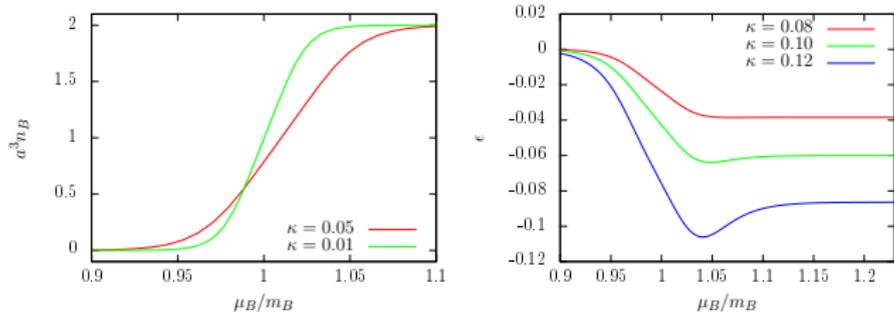
# “Solution” of the sign problem: Numerical and analytic investigations of the effective theory



- effective theory inherits only mild version of sign problem
- solution 1: standard MC simulations and reweighting
- solution 2: complex Langevin algorithm
- correctness criteria checked, consistent results
- solution 3: perturbative expansion in effective couplings

## Low temperature limit in the heavy dense regime

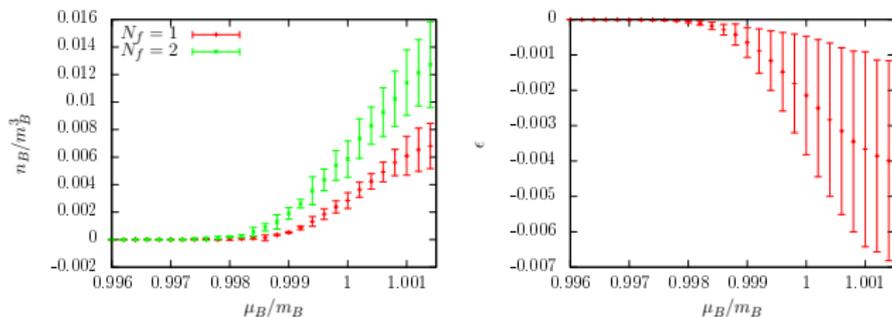
- low temperature:  $N_T$  large; heavy:  $\kappa \ll 1$ ; dense:  $c \approx 1$ ;  $\bar{c} \approx 0$
- ⇒ dominated by short range quark line interactions
- physics of the nuclear liquid-gas transition



### Features:

- $n_B$ : silver blaze and saturation by Pauli exclusion principle
- nuclear binding energy  $\epsilon = \frac{e - n_B m_B}{n_B m_B}$ : negative due to attractive quark interaction

## Continuum limit in the heavy dense low temperature regime



- continuum limit  $a \rightarrow 0$  at fixed  $\frac{m_B}{T}$  and  $T = \frac{1}{aN_\tau}$  requires larger value of  $\kappa$
  - truncation error: difference of  $\kappa^2$  and  $\kappa^4$  results
  - main features persist in the continuum limit
- ⇒ higher orders in hopping expansion of particular importance for the continuum limit

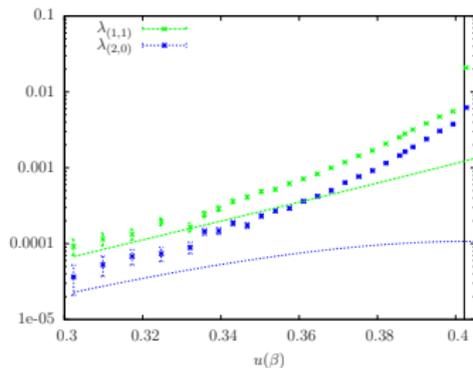
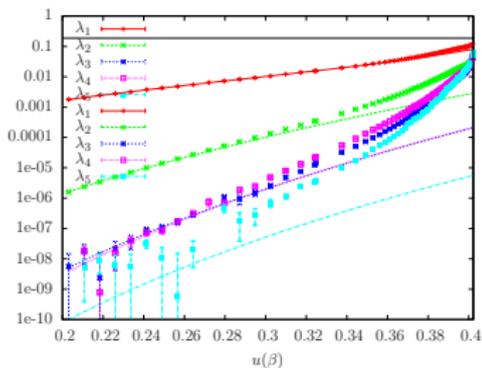
## Numerical determination of effective couplings

- strong coupling motivated form of effective action

$$e^{-S_{\text{eff}}} = \prod_{x, i=1, \dots, 3} \left( 1 + \lambda_1 (L(x)L(x + \hat{i})^\dagger + L(x)^\dagger L(x + \hat{i})) \right) \prod_{[x, y]} \left( 1 + \lambda_2 (L(x)L(y)^\dagger + L(x)^\dagger L(y)) \right) \dots$$

- include in the same way: long range interactions, different representations of Polyakov loops, n-point interactions
- assumptions: effective couplings  $\ll 1$ , long range interactions and higher representations suppressed
- coupling constants of the effective theory can be approximately related to correlators of Polyakov loops
- allows to estimate the couplings and check the strong coupling results

## Numerical results for the effective couplings

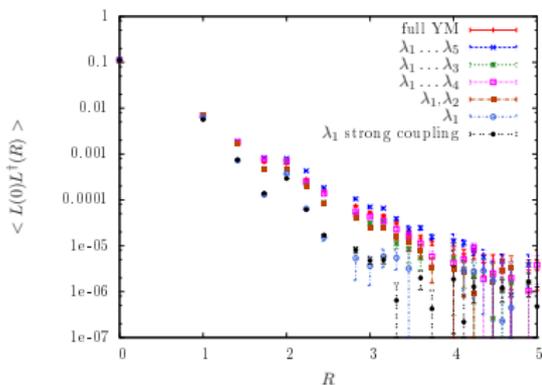


- next neighbor interaction: strong coupling result good approximation in wide parameter range
- long range interaction: strong coupling result less precise
- interactions of higher representations, higher n-point terms same size as long range interactions of fundamental Polyakov loop

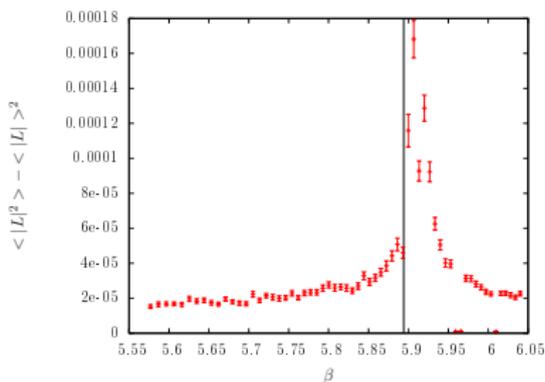
# Relevance of the difference to strong coupling approach

Depends on observable:

- correlators of fundamental Polyakov loops: significant influence of long range two point fund. interactions
- phase transition: dominated by short range interactions
- thermodynamics: higher rep. as important as long range two point fund. interactions



$\beta = 5.4; N_\tau = 4$



$N_\tau = 6; \text{two coupling model}$

## Conclusions

- effective Polyakov loop theory allows to circumvent the sign problem by analytic and numerical methods
- derivation of effective theory: improved strong coupling and hopping parameter expansion
- advantage of this approach: functional dependence of effective couplings on the bare parameters
- interesting application: heavy quark nuclear liquid-gas transition
- approach can be improved/checked using the effective couplings obtained in numerical simulations of the full theory
- implications of the differences to strong coupling approach depend on the observables
- If the small corrections of long range interactions of fund. loops are considered important for thermodynamics other, more involved, terms should be considered as well.